Adomian Decomposition Method For solving Fractional Differential Equations

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Abstract: The Adomian decomposition method (ADM) is a non numerical method for solving a wide variety of functional equations and usually gets the solution in a series form.

System of fractional partial differential equation which has numerous applications in many fields of science is considered. Adomian decomposition method, a novel method is used to solve these type of equations. The solutions are derived in convergent series form which shows the effectiveness of the method for solving a wide variety of fractional differential equations.

Keywords: Adomian decomposition method, Fractional partial differential equations, System of differential equations, initial value problems.

Introduction

In recent years, considerable interest in fractional differential equations has been stimulated due to their numerous applications in many fields of science and engineering[9,10]. One of these applications, Adomian decomposition method (ADM)introduced by provides Adomian (1980), an effective procedure for finding explicit and numerical solutions of a wider and general class of differential systems representing real physical problems [2,3]. This method efficiently works for initial value or boundary value problems, for linear or nonlinear, ordinary or partial differential equations[5,6], and even for stochastic systems as well. Moreover, no linearization or perturbation is required in this method.

2. PRELIMINARIES AND NOTATIONS: If

f(t) is continous on an interval [a, b] and $0 < a \le 1$

, then the operator $I_{0^+}^{\alpha}$, defined by

$$I_{0^{+}}^{\alpha} f(t) = \frac{1}{\Gamma(\alpha)} \int_{0}^{t} \frac{f(t)}{(t-s)^{1-\alpha}} ds$$
 (1)

is called the Riemann- Liouville fractional integral operator of order α . Here $\Gamma(.)$ is the Gamma function.

The Caputo time fractional derivative of order $\alpha > 0$ [16], is defined as: $D_t^{\alpha} u(x,t) = \frac{\partial^{\alpha} u(x,t)}{\partial t^{\alpha}} =$ $\left\{ \frac{1}{\Gamma(m-\alpha)} \int_0^t (t-\tau)^{m-\alpha-1} \frac{\partial^m u(x,t)}{\partial t^m}, (m-1 < \alpha < m) \\ \frac{\partial^m u(x,\tau)}{\partial t^m}, (m = \alpha \in N) \right\}$

Let $\rho, q \ge 0, f(t) \in L_1[0,T]$. Then $I_0^{\rho} + I_0^{q} + f(t) = I_{0^+}^{\rho+q} f(t) = I_{0^+}^{q} I_{0^+}^{\rho} f(t)$ is satisfied almost everywhere on [0,T]. Moreover, if $f(t) \in L_1[0,T]$, then the above equation is true for all $t \in [0,T]$.

Lemma2 (See [1])



If

$$\alpha > 0, f(t) \in L_1[0,T],$$
 then

 $c_{D_{0}^{\alpha}} I_{0}^{\alpha} f(t) = f(t)$, for all $t \in [0,T]$

Basic idea of Adomian Decomposition Method

Consider the differential equation.

Lu + Ru + Nu = g(3)

Where

L – Highest order derivative and easily invertible.

R – Linear differential operator of order less than L.

g – Source term.

Nu – Represents the nonlinear terms. The function, u(t) is assumed to be bounded for all $t \in I=[0,T]$ and the nonlinear term Nu satisfies Lipschitz condition i.e. $|Nu-Nv| \leq L_1 |u-v|$. Where

 L_1 is a positive constant.

Because *L* is invertible, we get

 $u = \Phi + L^{-1}g - L^{-1}Ru - L^{-1}Nu \quad (4)$

Where ϕ is the integration constant and satisfies $L\phi = 0$ and

$$L^{-1}(.) = \int_0^t (.) dt$$

The unkown function u is given by the infinite series.

$$u(x,t) = \sum_{n=0}^{\infty} u_n(x,t)$$
 (5)

And the nonlinear term *Nu*will be decomposed by the infinite series of Adomian polynomials

$$Nu = \sum_{n=0}^{\infty} A_n (u_0, u_1, ..., u_n)$$
(6)

Where A_n is Adomian polynomial calculated by using the formula

$$An = \frac{1}{n!} \left[\frac{d^n}{d\lambda^n} N\left(\nu(\lambda)\right) \right]_{\lambda=0} ,$$

$$n = 0, 1, 2, \dots$$

Where

$$\nu(\lambda) = \sum_{n=0}^{\infty} \lambda^n u_n$$

Substituting the decomposition series Eq. (5) and Eq. (6) into Eq. (4), gives

$$\begin{split} & \sum_{n=0}^{\infty} u_n(x,t) = \\ & \varphi + L^{-1}g - L^{-1}R(\sum_{n=0}^{\infty} u_n(x,t)) - \\ & \sum_{n=0}^{\infty} A_n(u_0, u_1, u_2, \dots, u_n) \\ & (7) \end{split}$$

From the above equation, we observe that

$$u_{0} = \varphi + L^{-1}g$$

$$u_{1} = -L^{-1}(Ru_{0}) - L^{-1}(A_{0})$$

$$u_{2} = -L^{-1}(Ru_{1}) - L^{-1}(A_{1})$$

:

$$u_{n+1} = -L^{-1}(Ru_n) - L^{-1}(A_n)$$
 , $n \ge 0$

Where φ is the initial condition.

Hence all terms of u are calculated and the general solution obtained according to ADM as $u(x,t) = \sum_{n=0}^{\infty} u_n(x,t)$

The convergence of this series has been proved [8].

Now, we apply Adomian decomposition method to derive the solution of fractional partial differential equations. We solve five examples by Adomian Decomposition Method.

Firstly, we apply the Adomian decomposition method to obtain approximate solutions of IVPs for fractional BBM-Burger's equation with $\mathcal{E} = 1$

Example 1.

Consider now the following equation:

$$\frac{\partial^{\alpha} u(x,t)}{\partial t^{\alpha}} = \left(L_{xx} u(x,t) \right) - \left(u(x,t) L_{x} u(x,t) \right)$$

Where $L_{xx} = \frac{\partial^{2}}{\partial x^{2}}$, $L_{x} = \frac{\partial}{\partial x}$

With the initial condition

 $u(x,0) = \sin x$, $(x,t) \in [0,1] \times (0,T]$

And the fractional differential operator $\frac{\partial^{\alpha}}{\partial t^{\alpha}}$ defined by (2).Let J^{α} be the inverse of the operator $\frac{\partial^{\alpha}}{\partial t^{\alpha}}$, now applying J^{α} to the both sides

of (10), we get

$$u(x,t) = \Phi + J^{\alpha}(L_{xx}u) - J^{\alpha}(Nu)$$

where $Nu = u L_x u$

In order to solve our problem,we must generalize these Adomian polynomials in follows.

$$A_{n} = \frac{1}{n!} \left[\frac{d^{n}}{d\lambda^{n}} N\left(\sum_{i=0}^{\infty} \lambda^{i} u_{i} \right) \left(\frac{\partial}{\partial x} \sum_{i=0}^{\infty} \lambda^{i} u_{i} \right) \right]_{\lambda=0}, n \ge 0$$

The first few terms of the Adomian polynomials are derived as follows

$$\begin{split} A_0 &= u_0 \frac{\partial u_0}{\partial x} \\ A_1 &= \frac{1}{1!} \left[\frac{d}{d\lambda} \left[(u_0 + \lambda u_1) \left(\frac{\partial u_0}{\partial x} + \lambda \frac{\partial u_1}{\partial x} \right) \right] \right]_{\lambda=0} \\ &= u_0 \frac{\partial u_1}{\partial x} + u_1 \frac{\partial u_0}{\partial x_1} \\ A_2 &= \frac{1}{2!} \left[\frac{d^2}{d\lambda^2} \left[(u_0 + \lambda u_1 + \lambda^2 u^2) \left(\frac{\partial u_0}{\partial x} + \lambda \frac{\partial u_1}{\partial x} + \lambda^2 \frac{\partial u_2}{\partial x} \right) \right] \right]_{\lambda=0} \end{split}$$

 $= u_0 \frac{\partial u_2}{\partial x} + u_1 \frac{\partial u_1}{\partial x} + u_2 \frac{\partial u_0}{\partial x}$ Similarly

$$A_{3} = u_{0}\frac{\partial u_{3}}{\partial x} + u_{1}\frac{\partial u_{2}}{\partial x} + u_{2}\frac{\partial u_{1}}{\partial x} + u_{3}\frac{\partial u_{0}}{\partial x}$$

:

And so on

From (7), we get

$$u(x,t) = \varphi + J^{\alpha}(\sum_{n=0}^{\infty} (L_{xx}u_n)) - - J^{\alpha}(\sum_{n=0}^{\infty} (A_n))$$

$$u_{0} = \varphi = u(x, 0)$$

$$u_{1} = J^{\alpha}(L_{xx}u_{0}) - J^{\alpha}(A_{0})$$

$$u_{2} = J^{\alpha}(L_{xx}u_{1}) - J^{\alpha}(A_{1})$$

$$\vdots$$

$$u_{n+1} = J^{\alpha}(L_{xx}u_{n}) - J^{\alpha}(A_{n})$$

By substituting the values of u_0, u_1, \ldots from above,

we get the solution of the IVP

$$\begin{aligned} u(x,t) &= u_0 + u_1 + \dots + u_n + \dots \\ u_0 &= u(x,0) = f(x) = \sin x \\ u_1 &= J^{\alpha} (L_{xx} u_0) - J^{\alpha} (A_0) \\ &= \frac{1}{\Gamma(\alpha)} \int_0^t \frac{1}{(t-\theta)^{1-\alpha}} f''(x) d\theta - \frac{1}{\Gamma(\alpha)} \int_0^t \frac{1}{(t-\theta)^{1-\alpha}} [f(x) \\ u_1 &= \frac{f_1(x)}{\Gamma(\alpha+1)} t^{\alpha} \\ \text{Where } f_1(x) &= -f''(x) + f(x)f'(x) \\ &= \sin(x)(1+\cos x). \\ u_2 &= J^{\alpha} (L_{xx} u_1) - J^{\alpha} (A_1) \\ &= f_2(x) \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} \end{aligned}$$



Where
$$f_2(x) = f_1''(x) - f(x)f_1'(x) - f_1(x)f'(x)$$

= $\sin^3(x) + [-1 - 5\cos x - \cos^2 x - (1 + \cos x)\cos x]\sin x$

$$\begin{split} u_3 &= J^{\alpha}(A_2) = \frac{1}{\Gamma(\alpha)} \int_0^t \frac{1}{(t-\theta)^{1-\alpha}} \Big[u_0 \frac{\partial u_2}{\partial x} + u_1 \frac{\partial u_1}{\partial x} + u_2 \frac{\partial u_0}{\partial x} \Big] d\theta \\ u_3 &= J^{\alpha}(L_{xx}u_2) - J^{\alpha}(A_2) \\ &= f_3(x) \frac{t^{3\alpha}}{\Gamma(3\alpha+1)} \end{split}$$

Similarly

$$u_4 = \frac{f_4(x)}{\Gamma(4\alpha + 1)} t^{4\alpha}$$

 $u_n = \frac{f_n(x)}{\Gamma(n\alpha + 1)} t^{n\alpha}$

The solution of the equation in series is given by

$$\begin{split} u(x,t) &= f(x) + \frac{t^{\alpha}}{\Gamma(\alpha+1)} f_1(x) + \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} f_2(x) + \frac{t^{3\alpha}}{\Gamma(3\alpha+1)} f_3(x) \dots + \dots + \frac{t^{n\alpha}}{\Gamma(n\alpha+1)} f_n(x) + \dots \\ &= \sum_{n=0}^{\infty} \frac{t^{n\alpha}}{\Gamma(n\alpha+1)} f_n(x) \quad , \text{ Where } f_0(x) \text{ is an initial condition.} \end{split}$$

Next, we will solve a more general system of nonlinear fractional differential equations.

Example 2.

Consider the following of nonlinear fractional differential equation:

$$\frac{\partial^{\alpha} u}{\partial t^{\alpha}} - \frac{\partial^{3} u}{\partial x^{3}} + \frac{\partial^{2} u}{\partial x^{2}} - \frac{\partial u}{\partial x} + u \frac{\partial u}{\partial x} = 0$$

(x, t) $\in \Omega \times (0, T]$ and $0 < \alpha \le 1$
With the initial condition $u(x, 0) = f(x)$

Where f(x) = cosx

$$\frac{\partial^{\alpha} u}{\partial t^{\alpha}} = -u \frac{\partial u}{\partial x} + \frac{\partial^{3} u}{\partial x^{3}} - \frac{\partial^{2} u}{\partial x^{2}} + \frac{\partial u}{\partial x}$$
$$= -u(x,t) L_{x} u(x,t) + L_{xxx} u(x,t) - L_{xx} u(x,t) + L_{x} u(x,t)$$
Where $L_{xxx} = \frac{\partial^{3}}{\partial x^{3}}$, $L_{xx} = \frac{\partial^{2}}{\partial x^{2}}$, $L_{x} = \frac{\partial}{\partial x}$

And the fractional D_t^{α} defined as Eq. (2) we know

that \int_{t}^{α} is inverse of the operator D_{t}^{α}

Now applying \int^{α} to the both side of the given

Eq., we obtain.

$$\begin{split} u(x,t) &= \varphi - J^{\alpha}(Nu) + J^{\alpha}(L_{xxx}u) - J^{\alpha}(L_{xx}u) + J^{\alpha}(L_{x}u) \end{split}$$
 Where $Nu = u \frac{\partial u}{\partial x}$

The first few terms of the Adomian polynomials are given by:

$$A_{0} = u_{0} \frac{\partial u_{0}}{\partial x}$$

$$A_{1} = u_{0} \frac{\partial u_{1}}{\partial x} + u_{1} \frac{\partial u_{0}}{\partial x}$$

$$A_{2} = u_{0} \frac{\partial u_{2}}{\partial x} + u_{1} \frac{\partial u}{\partial x} + u_{2} \frac{\partial u_{0}}{\partial x}$$

$$\vdots$$

And so on

From (7), we get

$$\sum_{n=0}^{\infty} u_n(x,t) = \varphi - J^{\alpha}(\sum_{n=0}^{\infty} A_n(u_0, u_1, \dots, u_n)) + J^{\alpha}(\sum_{n=0}^{\infty} L_{xxx}u_n) - J^{\alpha}(\sum_{n=0}^{\infty} L_{xx}u_n) + J^{\alpha}(\sum_{n=0}^{\infty} L_xu_n)$$

It is clear that:

$$\begin{split} u_1 &= -J^{\alpha}A_0 + J^{\alpha}L_{xxx}u_0 - J^{\alpha}L_{xx}u_0 + J^{\alpha}L_xu_0, \\ u_2 &= -J^{\alpha}A_1 + J^{\alpha}L_{xxx}u_1 - J^{\alpha}L_{xx}u_1 + J^{\alpha}L_x \\ \vdots \end{split}$$

$$\begin{split} u_{n+1} &= -J^{\alpha}A_n + J^{\alpha}L_{xxx}u_n - J^{\alpha}L_{xx} + J^{\alpha}L_xu_n. \\ \text{By substituting the value of } u_0 \ , u_1, \ldots \end{split}$$

From Eq. (20), we get the solution of the IVP



$$u(x,t) = u_0 + u_1 + u_2 + \dots + u_n + \dots$$

$$u_0 = u(x,0) = f(x) = \cos x \ u_1 = -J^{\alpha} A_0 + J^{\alpha} (L_{xxx} u_0) - J^{\alpha} (L_{xx} u_0) + J^{\alpha} (L_x u_0)$$

$$= \frac{f(x)f'(x) - f'''(x) + f''(x) - f'(x)}{\Gamma(\alpha + 1)} t^{\alpha}$$

$$\therefore \ u_1 = f_1(x) \frac{t^{\alpha}}{\Gamma(\alpha + 1)}$$

Where

$$f_{1}(x) = f(x)f'(x) - f'''(x) + f''(x) - f'(x)$$

= $(-\sin x - 1)\cos x - 2\sin x$
 $u_{2} = -J^{\alpha}(A_{1}) + J^{\alpha}(L_{xxx}u_{1}) - J^{\alpha}(L_{xx}u_{1}) + J^{\alpha}(L_{x}u_{1})$
= $\frac{f_{2}(x)}{\Gamma(2\alpha + 1)}t^{2\alpha}$

Where

$$\begin{split} f_2(x) &= -[f(x) + f_1'(x) + f_1(x)f'(x) - f_1'''(x) + \\ f_1''(x) - f_1'(x)] \end{split}$$

$$=\cos^{3}x + 3\cos^{2}x + \left((-\sin x - 1)\sin x - 2\sin x - 1\right)\cos x - 5\sin^{2}x - 2\sin x,$$

similarly

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 $u_3 = \frac{f_3(x)}{\Gamma(3\alpha+1)} t^{3\alpha}$

$$\begin{split} u_n &= \frac{f_n(x)}{\Gamma(n\alpha + 1)} t^{n\alpha} \\ u(x,t) &= f(x) + \frac{t^{\alpha}}{\Gamma(\alpha + 1)} f_1(x) + \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} f_2(x) + \dots + \frac{t^{n\alpha}}{\Gamma(n\alpha + 1)} f_n(x) + \dots \\ &= \sum_{n=0}^{\infty} \frac{t^{n\alpha}}{\Gamma(n\alpha + 1)} f_n(x) \\ \text{Where} \quad f_0(x) \text{is an initial condition.} \end{split}$$

In the following example, we try to find the solve of another nonlinear fractional equation

Example 3.

Consider the following of nonlinear fractional equation:-

$$D_t^{\alpha} u + D_x^2 u - D_x u + u^2 = 0$$

$$0 < x \le 1, 0 \le t \le 1, \quad 0 < \alpha \le 1$$

With initial condition

$$u(x, 0) = \varphi = f(x) = x^2,$$

$$(x,t) \in \Omega \times [0,T]$$

Note that here $\Omega = (0, 1)$

The standard form of the fractional equation an operator form is

$$D_t^{\alpha} u = -[u(x,t)]^2 - L_{xx}u(x,t) + L_x u(x,t)$$

Where $L_{xx} = \frac{\partial^2}{\partial x^2}$, $L_x = \frac{\partial}{\partial x}$ and the fractional differential operator D_t^{α} defined in (2) respectively.

 \int_{t}^{α} is the inverse of D_{t}^{α}

Now applying J^{α} to the both side of our Eq. we obtain

$$u(x,t) = \Phi - J^{\alpha}(Nu) - J^{\alpha}(L_{xx}u(x,t)) + J^{\alpha}(L_{x}u(x,t))$$

Where $Nu = u^{2}$, according to the decomposition
method, we assume series solution for the
unknown function $u(x,t)$ in the form

$$u(x,t) = \sum_{n=0}^{\infty} u_n(x,t)$$

In order to solve our problem, we must generalize these Adomian polynomials as follows:

$$An = \frac{1}{n!} \left[\frac{d^n}{d\lambda^n} \left(\sum_{i=0}^{\infty} \lambda^i u_i \right) \left(\sum_{i=0}^{\infty} \lambda^i u_i \right) \right]_{\lambda=0}, \quad n = 0, 1, \dots$$

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$$\begin{split} A_0 &= u_0 . u_0 = u_0^2 \\ A_1 &= \frac{d}{d\lambda} [(u_0 + \lambda u_1) . (u_0 + \lambda u_1)]_{\lambda = 0} \\ &= u_0 u_1 + u_1 u_0 = 2u_0 u_1 \\ A_2 &= \frac{1}{2!} \frac{d^2}{d\lambda^2} [(u_0 + \lambda u_1 + \lambda^2 u_2) (u_0 + \lambda u_1 + \lambda^2 u_2)] \\ &= 2u_0 u_2 + u_1^2 \\ A_3 &= \frac{1}{3!} \frac{d^3}{d\lambda^3} [(u_0 + \lambda u_1 + \lambda^2 u_2 + \lambda^3 u_3) (u_0 + \lambda u_1 + \lambda^2 u_2 + \lambda^3 u_3)]_{\lambda = 0} \\ &= 2u_0 u_3 + 2u_1 u_2 \\ \vdots \end{split}$$

$$\therefore u_{1} = x^{4} - 2x + 2 \left[\frac{t^{\alpha}}{\Gamma(\alpha + 1)} \right]$$

$$u_{2} = -J^{\alpha}(A_{1}) - J^{\alpha}(L_{xx}u_{1}) + J^{\alpha}(L_{x}u_{1})$$

$$= \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} f_{2}(x)$$
Where $f_{2}(x) = -[2f(x)f_{1}(x) + f_{1}''(x) - f_{1}'(x)]$

$$= -2x^{6} + 8x^{2} + 2$$

$$u_{3} = -J^{\alpha}(A_{2}) - J^{\alpha}(L_{xx}u_{2}) + J^{\alpha}(L_{x}u_{2})$$

$$= \frac{t^{3\alpha}}{\Gamma(3\alpha + 1)} f_{3}(x)$$

Where

$$f_3(x) = -2f(x)f_2(x) - 2f_1^2(x) - f_2''(x) + f_2'(x)$$

= $3x^8 - 8x^5 + 40x^4 - 8x^2 + 24x - 20$

Similarly

$$u_4 = \frac{t^{4\alpha}}{\Gamma(4\alpha + 1)} f_4(x)$$

$$u_n = \frac{t^{n\alpha}}{\Gamma(n\alpha+1)} f_n(x)$$

The solution of the considered IVP is given by

$$\begin{split} u(x,t) &= f(x) + \frac{t^{\alpha}}{\Gamma(\alpha+1)} f_1(x) + \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} f_2(x) + \cdots \\ &= \sum_{n=0}^{\infty} \frac{t^{n\alpha}}{\Gamma(n\alpha+1)} f_n(x) \end{split}$$

Where $f_0(x) = u(x, 0)$, is initial condition.

Also, in the next example, we solve a system of nonlinear equations with fractional orders **Example 4.**

Consider the system of initial value problem (IVP) of fractional equations

From (7), we get

$$\sum_{n=0}^{\infty} u_n(x,t) = \varphi - J^{\alpha}(\sum_{n=0}^{\infty} A_n(u_0, u_1, \dots, u_n)) - J^{\alpha}(\sum_{n=0}^{\infty} (L_{xx}u_n)) + J^{\alpha}(\sum_{n=0}^{\infty} (L_xu_n))$$

From the above equation, we observe that

$$\begin{split} u_0 &= u(x,0) = f(x) \\ u_1 &= -J^{\alpha}(A_0) - J^{\alpha}(L_{xx}u_0) + J^{\alpha}(L_xu_0) \\ u_2 &= -J^{\alpha}(A_1) - J^{\alpha}(L_{xx}u_1) + J^{\alpha}(L_xu_1) \\ \vdots \end{split}$$

$$u_{n+1} = -J^{\alpha}(A_n) - J^{\alpha}(L_{xx}u_n) + J^{\alpha}(L_xu_n)$$

By substituting the values of u_0 , u_1 , ... ,

we get solutions of the IVP

$$u = u_0 + u_1 + u_2 + \dots + u_n + \dots$$

$$u_0 = u(x, 0) = f(x) = x^2$$

$$u_1 = -J^{\alpha}(A_0) - J^{\alpha}(L_{xx}u_0) + J^{\alpha}(L_x(u_0))$$

$$= \frac{t^{\alpha}}{\Gamma(\alpha + 1)} f_1(x)$$

Where

$$f_1(x) = (f(x))^2 + f''(x) - f'(x) = x^4 - 2x + 2$$

 $D_t^{\alpha} u = u D_x u + v D_y u$

$$D_t^{\alpha} v = u D_x v + v D_y v$$

Where $0 < \alpha \leq 1$ and $(x, t) \in \Omega_x$ (0, T], and with

the initial condition.

u(x, y, 0) = f(x, y)

$$v(x,y,0) = g(x,y)$$
, $x,y \in \Omega$

Not that $\Omega = (0, 1)$

The above system can be written in the equivalent form

$$\frac{\partial^{\alpha} u}{\partial t^{\alpha}} = u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y}$$

$$\frac{\partial^{\alpha} v}{\partial t^{\alpha}} = u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y}$$

$$N_{1}(u,v) = u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y}$$

$$N_{2}(u,v) = u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y}$$

$$Lu = N_{1}(u,v)$$

$$Lv = N_{2}(u,v)$$

Applying $L^{-1}(.) = J^{\alpha}$ to both sides of Eq.(28) yields.

 $u(x, y, t) = \phi + J^{\alpha} N_1(u, v)$ $v(x, y, t) = \phi + J^{\alpha} N_2(u, v)$

Where the nonlinear operator $N_1(u,v)$ and $N_2(u,v)$ are them written in the decomposition form

$$N_{1}(u,v) = \sum_{n=0}^{\infty} A_{n} (u_{0}, u_{1}, \dots, u_{n})$$
$$N_{2}(u,v) = \sum_{n=0}^{\infty} B_{n} (v_{0}, v_{1}, \dots, v_{n})$$

Where A_n and B_n are the Adomian polynomials of the following form

$$A_n = \frac{1}{n!} \left[\frac{d^n}{d\lambda^n} N_1(u, v) \left(\sum_{i=0}^{\infty} \lambda^i u_i \right) \right]_{\lambda=0}$$

$$B_n = \frac{1}{n!} \left[\frac{d^n}{d\lambda^n} N_2(u, v) \left(\sum_{i=0}^{\infty} \lambda^i v_i \right) \right]_{\lambda=0}, n = 0, 1, \dots$$

generalize these Adomian polynomials in follows.

$$\begin{split} A_n &= \frac{1}{n!} \Big[\frac{d^n}{d\lambda^n} \left(\sum_{i=0}^{\infty} \lambda^i u_i \right) \left(\frac{\partial}{\partial x} \sum_{i=0}^{\infty} \lambda^i u_i \right) + \\ & \left(\sum_{i=0}^{\infty} \lambda^i v_i \right) \left(\frac{\partial}{\partial y} \sum_{i=0}^{\infty} \lambda^i u_i \right) \Big]_{\lambda=0} \quad , \end{split}$$

$$n = 0, 1, ...$$

$$A_{0} = u_{0} \frac{\partial u_{0}}{\partial x} + v_{0} \frac{\partial u_{0}}{\partial y}$$

$$\therefore A_{1} = u_{0} \frac{\partial u_{1}}{\partial x} + v_{0} \frac{\partial u_{1}}{\partial y} + u_{1} \frac{\partial u_{0}}{\partial x} + v_{1} \frac{\partial u_{0}}{\partial y}$$

$$\therefore A_{2} = u_{0} \frac{\partial u_{2}}{\partial x} + v_{0} \frac{\partial u_{2}}{\partial y} + u_{1} \frac{\partial u_{1}}{\partial x} + v_{1} \frac{\partial u_{1}}{\partial y} + u_{2} \frac{\partial u_{0}}{\partial x} + v_{2} \frac{\partial u_{0}}{\partial y}$$

Similarly

$$: A_3 = u_0 \frac{\partial u_3}{\partial x} + v_0 \frac{\partial u_3}{\partial y} + u_1 \frac{\partial u_2}{\partial x} + v_1 \frac{\partial u_2}{\partial y} + u_2 \frac{\partial u_1}{\partial x} + v_2 \frac{\partial u_1}{\partial y} + u_3 \frac{\partial u_0}{\partial x} + v_3 \frac{\partial u_0}{\partial y}$$

And so on

Now, to calculated our problem we must generalize these Adomian polynomials in follows



$$B_{n} = \frac{1}{n!} \left[\frac{d^{n}}{d\lambda^{n}} \left(\sum_{i=0}^{\infty} \lambda^{i} u_{i} \right) \left(\frac{\partial}{\partial x} \sum_{i=0}^{\infty} \lambda^{i} v_{i} \right) + \left(\sum_{i=0}^{\infty} \lambda^{i} v_{i} \right) \left(\frac{\partial}{\partial y} \sum_{i=0}^{\infty} \lambda^{i} v_{i} \right) \right]_{\lambda=0}, \quad u_{2} = J^{\alpha} A_{1}$$

$$v_{2} = J^{\alpha} B_{1}$$

$$\vdots$$

$$n = 0, 1, 2, ...$$

$$u_{n+1} = J^{\alpha} A_{n}$$

$$v_{n+1} = J^{\alpha} B_{n}$$

$$B_{0} = u_{0} \frac{\partial u_{0}}{\partial x} + v_{0} \frac{\partial v_{0}}{\partial y}$$

$$B_{1} = u_{0} \frac{\partial v_{1}}{\partial x} + v_{0} \frac{\partial v_{1}}{\partial y} + u_{1} \frac{\partial v_{0}}{\partial x} + v_{1} \frac{\partial v_{0}}{\partial y}$$

$$B_{2} = u_{0} \frac{\partial v_{2}}{\partial x} + v_{0} \frac{\partial v_{2}}{\partial y} + u_{1} \frac{\partial v_{1}}{\partial x} + v_{1} \frac{\partial v_{1}}{\partial y} + u_{2} \frac{\partial v_{0}}{\partial x} + v_{2} \frac{\partial v_{0}}{\partial y}$$

$$\therefore B_{3} = u_{0} \frac{\partial v_{3}}{\partial x} + v_{0} \frac{\partial v_{2}}{\partial y} + u_{1} \frac{\partial v_{2}}{\partial x} + v_{2} \frac{\partial v_{1}}{\partial y} + u_{2} \frac{\partial v_{0}}{\partial x} + v_{3} \frac{\partial v_{0}}{\partial y}$$

$$u_{1} = J^{\alpha} (A_{0})$$

And so on

From (7), we obtain

$$\sum_{n=0}^{\infty} u_n(x, y, t) = u(x, y, 0) + J^{\alpha}(\sum_{n=0}^{\infty} A_n(u_0, u_1, \dots, u_n))$$
$$\sum_{n=0}^{\infty} v_n(x, y, t) = v(x, y, 0) + J^{\alpha}(\sum_{n=0}^{\infty} B_n(u_0, u_1, \dots, u_n))$$

The associated decomposition is given by

$$u_{0} = u(x, y, 0), u_{n+1} = J^{\alpha}(N_{1}(u_{n}, v_{n}))$$

$$v_{0} = v(x, y, 0), v_{n+1} = J^{\alpha}(N_{2}(u_{n}, v_{n})), n = 0, 1, 2, ...$$

Then, According to the above equations we get

$$u_{0} = u(x, y, 0)$$

$$v_0 = v(x, y, 0)$$
$$u_1 = J^{\alpha} A_0$$

 $v_1 = J^{\alpha}B_0$

n ubstituting the values And v_0 , v_1 , ... we get a solution of the

$$u_{0} = u(x, y, 0) = f(x, y)$$

$$v_{0} = v(x, y, 0) = g(x, y)$$

$$u_{1} = J^{\alpha}(A_{0}) = \frac{1}{\Gamma(\alpha)} \int_{0}^{t} \frac{1}{(t-\theta)^{1-\alpha}} \left[u_{0} \frac{\partial u_{0}}{\partial x} + v_{0} \frac{\partial u_{0}}{\partial y} \right] d\theta$$

$$= f_{1}(x, y) \frac{t^{\alpha}}{\Gamma(\alpha+1)}$$

Where

$$\begin{split} f_1(x,y) &= -\left[f(x,y)\frac{\partial f(x,y)}{\partial x} + g(x,y)\frac{\partial f(x,y)}{\partial y}\right] \\ v_1 &= J^{\alpha}B_0 = \frac{1}{\Gamma(\alpha)}\int_0^t \frac{1}{(t-\theta)^{1-\alpha}}\left[u_0\frac{\partial v_0}{\partial x} + v_0\frac{\partial v_0}{\partial y}\right]d\theta \\ &= g_1(x,y)\frac{t^{\alpha}}{\Gamma(\alpha+1)} \end{split}$$

Where

$$g_1(x, y) = -\left[f(x, y)\frac{\partial g(x, y)}{\partial x} + g(x, y)\frac{\partial g(x, y)}{\partial y}\right]$$
$$u_2 = J^{\alpha}(A_1) = f_2(x)\frac{t^{2\alpha}}{\Gamma(2\alpha + 1)}$$

Where

$$f_2(x) = \left[f(x,y) \frac{\partial f_1(x,y)}{\partial x} + g(x,y) \frac{\partial f_1(x,y)}{\partial y} + f_1(x,y) \frac{f(x,y)}{\partial x} + g_1(x,y) \frac{\partial f(x,y)}{\partial y} \right]$$
$$v_2 = J^{\alpha} B_1 = g_2(x,y) \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)}$$

Where

$$g_{2}(x, y) = f(x, y) \frac{\partial g_{1}(x, y)}{\partial x} + g(x, y) \frac{\partial g_{1}(x, y)}{\partial y} + f_{1}(x, y) \frac{\partial g(x, y)}{\partial x} + g_{1}(x, y) \frac{\partial g(x, y)}{\partial y}$$

By induction, we have

$$u(x, y, t) = \sum_{n=0}^{\infty} u_n = u_0 + u_1 + u_2 + \dots + u_n + \dots$$

= $f(x, y) + f_1(x, y) \frac{t^{\alpha}}{\Gamma(\alpha + 1)} + f_2(x, y) \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} + u_1(x, y) \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} + \dots$
= $g(x, y) + g_1(x, y) \frac{t^{\alpha}}{\Gamma(\alpha + 1)} + g_2(x, y) \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} + \dots + g_n(x, y) \frac{t^{n\alpha}}{\Gamma(n\alpha + 1)} + \dots$

Finally, we apply the Adomian decomposition method to obtain approximate solutions of IVPs for fractional BBM-Burger's equation with $\mathcal{E} = 1$

Example 5.

Consider the initial value problem (IVP) for fractional **BBM_Burger's** equation of the form

$$\frac{\partial^{\alpha} u}{\partial t^{\alpha}} - \frac{\partial^2 u}{\partial x^2} + u \frac{\partial u}{\partial x} = 0$$

Where $0 < \alpha \leq 1$ and with initial condition

$$u(x,0) = \varphi = f(x) = \sin(x), x \in \Omega \times (0,T]$$

Note that here $\Omega = (0,1)$,the standard form of the fractional **BBM_Burger's** equation in an operator form

$$\frac{\partial^{\alpha} u}{\partial t^{\alpha}} = \frac{\partial^{2} u}{\partial x^{2}} - u \frac{\partial u}{\partial x}$$
$$\frac{\partial^{\alpha} u(x,t)}{\partial t^{\alpha}} = \left(L_{xx} u(x,t) \right) - \left(u(x,t) L_{x} u(x,t) \right)$$

Where $L_{xx} = \frac{\partial^2}{\partial x^2}$, $L_x = \frac{\partial}{\partial x}$

And the fractional differential operator $\frac{\partial^{\alpha}}{\partial t^{\alpha}}$ defined in equation (2), respectively we know that J^{α} which is invers of the operator $\frac{\partial^{\alpha}}{\partial t^{\alpha}}$, now applying J^{α} to the both sides of our Eq., we get

$$u(x,t) = \phi + J^{\alpha}(L_{xx}u) - J^{\alpha}(Nu)$$

+ ... Where $Nu_{f_n(x, y)} = \frac{u \frac{\partial w}{\partial x}^{\alpha}}{\Gamma(n\alpha + 1)} + \cdots$

In order to solve our problem we must generalize these Adomian polynomials as follows.

$$A_n = \frac{1}{n!} \frac{d^n}{d\lambda^n} \left[\left(\sum_{i=0}^n \lambda^i u_i \right) \frac{\partial}{\partial x} \left(\sum_{i=0}^n \lambda^i u_i \right) \right]_{\lambda=0}$$

$$n = 0, 1, 2, \dots$$

$$A_{0} = u_{0} \frac{\partial u_{0}}{\partial x}$$

$$A_{1} = u_{0} \frac{\partial u_{1}}{\partial x} + u_{1} \frac{\partial u_{0}}{\partial x}$$

$$A_{2} = u_{0} \frac{\partial u_{2}}{\partial x} + u_{1} \frac{\partial u_{1}}{\partial x} + u_{2} \frac{\partial u_{0}}{\partial x}$$

$$\vdots$$

And so on

Thus

$$u(x,t) = \varphi + J^{\alpha}(\sum_{n=0}^{\infty} (L_{xx}u_n)) - J^{\alpha}(\sum_{n=0}^{\infty} (A_n)),$$
$$u_{\alpha} = \varphi = u(x,0)$$

$$u_{1} = J^{\alpha}(L_{xx}u_{0}) - J^{\alpha}(A_{0})$$
$$u_{2} = J^{\alpha}(L_{xx}u_{1}) - J^{\alpha}(A_{1})$$
$$\vdots$$

$$u_{n+1} = J^{\alpha}(L_{xx}u_n) - J^{\alpha}(A_n)$$

Consequently

 $u(x,t) = u_0 + u_1 + u_2 + \dots + u_n + \dots$

 $u_0 = u(x, 0) = f(x) = \sin x$ $u_1 = J^{\alpha}(L_{xx}u_0) - J^{\alpha}(A_0) = \frac{f_1(x)}{\Gamma(\alpha + 1)}t^{\alpha}$ Where $f_1(x) = -f''(x) + f(x)f'(x)$ $= \sin x (1 + \cos x)$ $u_2 = J^{\alpha}(L_{xx}u_1) - J^{\alpha}(A_1) = \frac{f_2(x)}{\Gamma(\alpha + 1)}t^{\alpha}$

Where

$$\begin{aligned} f_2(x) &= \\ f_1''(x) - f(x)f_1'(x) - \\ f_1(x)f'(x) \end{aligned}$$

$$= \sin^{3} x + [-1 - 5\cos x - \cos^{2} x - (1 + \cos x)\cos x]$$
$$u_{3} = J^{\alpha}(L_{xx}u_{2}) - J^{\alpha}(A_{2}) = f_{3}(x)\frac{t^{3\alpha}}{\Gamma(3\alpha + 1)}$$

Similarly

$$\therefore u_3 = \frac{f_3(x)}{\Gamma(3\alpha+1)} t^{3\alpha}$$

÷

 $u_n = \frac{f_n(x)}{\Gamma(n\alpha + 1)} t^{n\alpha}$

The solution of the considered IVP is given by

$$\begin{split} u(x,t) &= f(x) + \frac{t^{\alpha}}{\Gamma(\alpha+1)} f_1(x) + \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} f_2(x) + \\ \frac{t^{3\alpha}}{\Gamma(3\alpha+1)} f_3(x) + \dots + \frac{t^{n\alpha}}{\Gamma(n\alpha+1)} f_n(x) + \dots \\ (38) \\ &= \sum_{n=0}^{\infty} \frac{t^{n\alpha}}{\Gamma(n\alpha+1)} f_n(x) , \end{split}$$

where $f_0(x)$ is an initial condition. (Comp. 7-23).

Conclusion

In this paper, we have applied the Adomian decomposition method for solving problems of nonlinear partial equations. We demonstrated that the decomposition procedure is quite efficient to determine the exact solutions. However, the method gives a simple powerful tool for obtaining the solutions without a need for large size of computations. It is also worth noting that the advantage of this method sometimes displays a fast convergence of the solutions. In addition, the numerical results which obtained by this method indicate a high degree of accuracy, also efficiency of the desired results.

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