# Adomian Decomposition Method For solving Fractional Differential Equations 

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#### Abstract

The Adomian decomposition method (ADM) is a non numerical method for solving a wide variety of functional equations and usually gets the solution in a series form. System of fractional partial differential equation which has numerous applications in many fields of science is considered. Adomian decomposition method, a novel method is used to solve these type of equations. The solutions are derived in convergent series form which shows the effectiveness of the method for solving a wide variety of fractional differential equations.


Keywords: Adomian decomposition method, Fractional partial differential equations, System of differential equations, initial value problems.

## Introduction

In recent years, considerable interest in fractional differential equations has been stimulated due to their numerous applications in many fields of science and engineering[9,10]. One of these applications, Adomian decomposition method (ADM)introduced by Adomian (1980), provides an effective procedure for finding explicit and numerical solutions of a wider and general class of differential systems representing real physical problems [2,3]. This method efficiently works for initial value or boundary value problems, for linear or nonlinear, ordinary or partial differential equations[5,6], and even for stochastic systems as well. Moreover, no linearization or perturbation is required in this method.

## 2. PRELIMINARIES AND NOTATIONS:

 If$f(t)$ is continous on an interval $[a, b]$ and $0<$ $\alpha \leq 1$
,then the operator $\mathrm{I}_{0^{+}}^{\alpha}$, defined by

$$
\begin{equation*}
\mathrm{I}_{0^{+}}^{\alpha} f(t)=\frac{1}{r(\alpha)} \int_{0}^{t} \frac{f(t)}{(t-s)^{1-\alpha}} d s \tag{1}
\end{equation*}
$$

is called the Riemann- Liouville fractional integral operator of order $\alpha$. Here $\Gamma($.$) is the$ Gamma function.

The Caputo time fractional derivative of order $\alpha>0$ [16], is defined as: $D_{t}^{\alpha} u(x, t)=\frac{\partial^{\alpha} u(x, t)}{\partial t^{\alpha}}=$
$\left\{\begin{array}{lr}\frac{1}{\Gamma(m-\alpha)} \int_{0}^{t}(t-\tau)^{m-\alpha-1} \frac{\partial^{m} u(x, t)}{\partial t^{m}},(m-1<\alpha<m) \\ \frac{\partial^{m} u(x, \tau)}{\partial t^{m}} & (m=\alpha \in N)\end{array}\right\}$

Lemma 1 (See [1])
Let $\quad \rho, q \geq 0, f(t) \in L_{1}[0, T]$. Then $\mathrm{I}_{0^{\mathrm{p}}+\mathrm{I}_{0^{+}}^{\mathrm{q}}} f(t)=\mathrm{I}_{0^{+}}^{\mathrm{p}+a} f(t)=\mathrm{I}_{0^{+}}^{a} \mathrm{I}_{0^{+}}^{\mathrm{p}} f(t) \quad$ is satisfied almost everywhere on $[0, T]$. Moreover, if $f(t) \in L_{1}[0, T]$, then the above equation is true for all $t \in[0, T]$.

Lemma2 (See [1])

If $\quad \alpha>0, f(t) \in L_{1}[0, T]$, then
$c_{D_{0^{+}}^{\alpha}} I_{0^{+}}^{\alpha} f(t)=f(t)$, for all $t \in[0, T]$
Basic idea of Adomian Decomposition

## Method

Consider the differential equation.
$L u+R u+N u=g$
Where
$L$ - Highest order derivative and easily invertible.
$R$ - Linear differential operator of order less than L ,
$g$-Source term.
Nu -Represents the nonlinear terms. The function, $u(t)$ is assumed to be bounded for all $\mathrm{t} \in \mathrm{I}=[0, \mathrm{~T}]$ and the nonlinear term Nu satisfies Lipschitz condition i.e. $|N u-N v| \leq L_{1}|u-v|$. Where $L_{1}$ is a positive constant.

Because $L$ is invertible, we get
$u=\phi+L^{-1} g-L^{-1} R u-L^{-1} N u$
Where $\phi$ is the integration constant and satisfies
$\mathrm{L} \phi=0$ and
$L^{-1}()=.\int_{0}^{t}() d$.
The unkown function $u$ is given by the infinite series,
$u(x, t)=\sum_{n=0}^{\infty} u_{n}(x, t)$
And the nonlinear term $N u$ will be decomposed by the infinite series of Adomian polynomials
$N u=\sum_{n=0}^{\infty} A_{n}\left(u_{0}, u_{1}, \ldots, u_{n}\right)$

Where $A_{n}$ is Adomian polynomial calculated by using the formula
$A n=\frac{1}{n!}\left[\frac{d^{n}}{d \lambda^{n}} N(v(\lambda))\right]_{\lambda=0}$,
$n=0,1,2, \ldots$
Where
$v(\lambda)=\sum_{n=0}^{\infty} \lambda^{n} u_{n}$
Substituting the decomposition series Eq. (5) and Eq. (6) into Eq. (4), gives

$$
\begin{aligned}
& \sum_{n=0}^{\infty} u_{n}(x, t)= \\
& \varphi+L^{-1} g-L^{-1} R\left(\sum_{n=0}^{\infty} u_{n}(x, t)\right)- \\
& \sum_{n=0}^{\infty} A_{n}\left(u_{0}, u_{1}, u_{2}, \ldots, u_{n}\right)
\end{aligned}
$$

(7)

From the above equation, we observe that
$u_{0}=\varphi+L^{-1} g$
$u_{1}=-L^{-1}\left(R u_{0}\right)-L^{-1}\left(A_{0}\right)$
$u_{2}=-L^{-1}\left(R u_{1}\right)-L^{-1}\left(A_{1}\right)$
:
$u_{n+1}=-L^{-1}\left(R u_{n}\right)-L^{-1}\left(A_{n}\right), \quad n \geq 0$

Where $\varphi$ is the initial condition.
Hence all terms of $u$ are calculated and the general solution obtained according to ADM as $u(x, t)=\sum_{n=0}^{\infty} u_{n}(x, t)$

The convergence of this series has been proved [8].

Now, we apply Adomian decomposition method to derive the solution of fractional partial differential equations.We solve five examples by Adomian Decomposition Method.

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Firstly, we apply the Adomian decomposition method to obtain approximate solutions of IVPs for fractional BBM-Burger's equation with $\varepsilon=1$

## Example 1.

Consider now the following equation:
$\frac{\partial^{\alpha} u(x, t)}{\partial t^{\alpha}}=\left(L_{x x} u(x, t)\right)-\left(u(x, t) L_{x} u(x, t)\right)$
Where $L_{x x}=\frac{\partial^{2}}{\partial x^{2}}, L_{x}=\frac{\partial}{\partial x}$
With the initial condition
$u(x, 0)=\sin x, \quad(x, t) \in[0,1] \times(0, T]$
And the fractional differential operator $\frac{\partial^{\alpha}}{\partial t^{\alpha}}$ defined by (2).Let $J^{\alpha}$ be the inverse of the operator $\frac{\partial^{\alpha}}{\partial t^{\alpha}}$, now applying $J^{\alpha}$ to the both sides of (10), we get
$u(x, t)=\phi+J^{\alpha}\left(L_{x x} u\right)-J^{\alpha}(N u)$
where $N u=u L_{x} u$
In order to solve our problem, we must generalize these Adomian polynomials in follows.
$A_{n}=\frac{1}{n!}\left[\frac{d^{n}}{d \lambda^{n}} N\left(\sum_{i=0}^{\infty} \lambda^{i} u_{i}\right)\left(\frac{\partial}{\partial x} \sum_{i=0}^{\infty} \lambda^{i} u_{i}\right)\right]_{\lambda=0}, n \geq 0$
The first few terms of the Adomian polynomials are derived as follows
$A_{0}=u_{0} \frac{\partial u_{0}}{\partial x}$
$A_{1}=\frac{1}{1!}\left[\frac{d}{d \lambda}\left[\left(u_{0}+\lambda u_{1}\right)\left(\frac{\partial u_{0}}{\partial x}+\lambda \frac{\partial u_{1}}{\partial x}\right)\right]\right]_{\lambda=0}$
$=u_{0} \frac{\partial u_{1}}{\partial x}+u_{1} \frac{\partial u_{0}}{\partial x_{1}}$
$A_{2}=\frac{1}{2!}\left[\frac{d^{2}}{d \lambda^{2}}\left[\left(u_{0}+\lambda u_{1}+\lambda^{2} u^{2}\right)\left(\frac{\partial u_{0}}{\partial x}+\lambda \frac{\partial u_{1}}{\partial x}+\lambda^{2} \frac{\partial u_{2}}{\partial x}\right)\right]\right]_{\lambda=0}$
$=u_{0} \frac{\partial u_{2}}{\partial x}+u_{1} \frac{\partial u_{1}}{\partial x}+u_{2} \frac{\partial u_{0}}{\partial x}$
Similarly
$A_{3}=u_{0} \frac{\partial u_{3}}{\partial x}+u_{1} \frac{\partial u_{2}}{\partial x}+u_{2} \frac{\partial u_{1}}{\partial x}+u_{3} \frac{\partial u_{0}}{\partial x}$
:
And so on
From (7), we get
$u(x, t)=$
$\varphi+J^{\alpha}\left(\sum_{n=0}^{\infty}\left(L_{x x} u_{n}\right)\right)-$
$-J^{\alpha}\left(\sum_{n=0}^{\infty}\left(A_{n}\right)\right)$
$u_{0}=\varphi=u(x, 0)$
$u_{1}=J^{\alpha}\left(L_{x x} u_{0}\right)-J^{\alpha}\left(A_{0}\right)$
$u_{2}=J^{\alpha}\left(L_{x x} u_{1}\right)-J^{\alpha}\left(A_{1}\right)$
:
$u_{n+1}=J^{\alpha}\left(L_{x x} u_{n}\right)-J^{\alpha}\left(A_{n}\right)$

By substituting the values of $u_{0}, u_{1}, \ldots$ from above, we get the solution of the IVP
$u(x, t)=u_{0}+u_{1}+\cdots+u_{n}+\cdots$
$u_{0}=u(x, 0)=f(x)=\sin x$
$u_{1}=J^{\alpha}\left(L_{x x} u_{0}\right)-J^{\alpha}\left(A_{0}\right)$
$=\frac{1}{\Gamma(\alpha)} \int_{0}^{t} \frac{1}{(t-\theta)^{1-\alpha}} f^{\prime \prime}(x) d \theta-\frac{1}{\Gamma(\alpha)} \int_{0}^{t} \frac{1}{(t-\theta)^{1-\alpha}}[f(x$
$u_{1}==\frac{f_{1}(x)}{\Gamma(\alpha+1)} t^{\alpha}$
Where $f_{1}(x)=-f^{\prime \prime}(x)+f(x) f^{\prime}(x)$
$=\sin (x)(1+\cos x)$.
$u_{2}=J^{\alpha}\left(L_{x x} u_{1}\right)-J^{\alpha}\left(A_{1}\right)$
$=f_{2}(x) \frac{t^{2 \alpha}}{\Gamma(2 \alpha+1)}$

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Where $f_{2}(x)=f_{1}^{\prime \prime}(x)-f(x) f_{1}^{\prime}(x)-f_{1}(x) f^{\prime}(x)$
$=\sin ^{3}(x)+\left[-1-5 \cos x-\cos ^{2} x-\right.$
$(1+\cos x) \cos x] \sin x$
$u_{3}=J^{\alpha}\left(A_{2}\right)=\frac{1}{\Gamma(\alpha)} \int_{0}^{t} \frac{1}{(t-\theta)^{1-\alpha}}\left[u_{0} \frac{\partial u_{2}}{\partial x}+u_{1} \frac{\partial u_{1}}{\partial x}+u_{2} \frac{\partial u_{0}}{\partial x}\right] d \theta$
$u_{3}=J^{\alpha}\left(L_{x x} u_{2}\right)-J^{\alpha}\left(A_{2}\right)$
$=f_{3}(x) \frac{t^{3 \alpha}}{\Gamma(3 \alpha+1)}$
Similarly
$u_{4}=\frac{f_{4}(x)}{\Gamma(4 \alpha+1)} t^{4 \alpha}$
$\vdots$
$u_{n}=\frac{f_{n}(x)}{\Gamma(n \alpha+1)} t^{n \alpha}$
The solution of the equation in series is given by
$u(x, t)=f(x)+\frac{t^{a}}{\Gamma(a+1)^{2}} f_{1}(x)+\frac{t^{2 n}}{\Gamma(2 a+1)^{2}} f_{2}(x)+\frac{t^{3 n}}{\Gamma(3 a+1)^{2}} f_{3}(x) \ldots+\cdots+\frac{t^{n a}}{\Gamma(n a+1)^{n n}} f_{n}(x)+\cdots$
$=\sum_{n=0}^{\infty} \frac{t^{n \alpha}}{\Gamma(n \alpha+1)} f_{n}(x) \quad$, Where $f_{0}(x)$ is an initial condition.

Next, we will solve a more general system of nonlinear fractional differential equations.

## Example 2.

Consider the following of nonlinear fractional differential equation:
$\frac{\partial^{\alpha} u}{\partial t^{\alpha}}-\frac{\partial^{3} u}{\partial x^{3}}+\frac{\partial^{2} u}{\partial x^{2}}-\frac{\partial u}{\partial x}+u \frac{\partial u}{\partial x}=0$
$(x, t) \in \Omega \times(0, T]$ and $0<\alpha \leq 1$
With the initial condition $u(x, 0)=f(x)$
Where $f(x)=\cos x$
$\frac{\partial^{\alpha} u}{\partial t^{\alpha}}=-u \frac{\partial u}{\partial x}+\frac{\partial^{3} u}{\partial x^{3}}-\frac{\partial^{2} u}{\partial x^{2}}+\frac{\partial u}{\partial x}$

$$
=-u(x, t) L_{x} u(x, t)+L_{x x x} u(x, t)-L_{x x} u(x, t)+L_{x}
$$

Where $L_{x x x}=\frac{\partial^{\mathrm{s}}}{\partial x^{\mathrm{s}}}, L_{x x x}=\frac{\partial^{2}}{\partial x^{2}}, L_{x}=\frac{\partial}{\partial x}$
And the fractional $D_{t}^{\alpha}$ defined as Eq. (2) we know that $I^{\alpha}$ is inverse of the operator $D_{t}^{\alpha}$

Now applying $J^{\alpha}$ to the both side of the given
Eq., we obtain.
$u(x, t)=\phi-J^{\alpha}(N u)+J^{\alpha}\left(L_{x x x} u\right)-J^{\alpha}\left(L_{x x} u\right)+J^{\alpha}\left(L_{x} u\right)$
Where $N u=u \frac{\partial u}{\partial x}$
The first few terms of the Adomian polynomials are given by:
$A_{0}=u_{0} \frac{\partial u_{0}}{\partial x}$
$A_{1}=u_{0} \frac{\partial u_{1}}{\partial x}+u_{1} \frac{\partial u_{0}}{\partial x}$
$A_{2}=u_{0} \frac{\partial u_{2}}{\partial x}+u_{1} \frac{\partial u}{\partial x}+u_{2} \frac{\partial u_{0}}{\partial x}$
:
And so on
From (7), we get
$\sum_{n=0}^{\infty} u_{n}(x, t)=\varphi-J^{a}\left(\sum_{n=0}^{\infty} A_{n}\left(u_{0}, u_{1}, \cdots, u_{n}\right)+J^{u}\left(\sum_{n=0}^{\infty} L_{x x x^{\prime}} u_{n}\right)-J^{u}\left(\sum_{n=0}^{\infty} L_{x u_{n}} u_{n}\right)+J^{u}\left(\sum_{n=0}^{\infty} L_{x} u_{n}\right)\right.$
It is clear that:
$u_{1}=-J^{\alpha} A_{0}+J^{\alpha} L_{x x x} u_{0}-J^{\alpha} L_{x x} u_{0}+J^{\alpha} L_{x} u_{0}$,
$u_{2}=-J^{\alpha} A_{1}+J^{\alpha} L_{x x x} u_{1}-J^{\alpha} L_{x x} u_{1}+J^{\alpha} L_{x}$
:
$u_{n+1}=-J^{\alpha} A_{n}+J^{\alpha} L_{x x x} u_{n}-J^{\alpha} L_{x x}+J^{\alpha} L_{x} u_{n}$.
By substituting the value of $u_{0}, u_{1}, \ldots$
From Eq. (20), we get the solution of the IVP

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$u(x, t)=u_{0}+u_{1}+u_{2}+\cdots+u_{n}+\cdots$
$u_{0}=u(x, 0)=f(x)=\cos x u_{1}=-J^{a} A_{0}+J^{a}\left(L_{x z x} u_{0}\right)-J^{a}\left(L_{x x} u_{0}\right)+J^{a}\left(L_{x} u_{0}\right)$

$$
=\frac{f(x) f^{\prime}(x)-f^{\prime \prime \prime}(x)+f^{\prime \prime}(x)-f^{\prime}(x)}{\Gamma(\alpha+1)} t^{\alpha}
$$

$\therefore u_{1}=f_{1}(x) \frac{t^{\alpha}}{\Gamma(\alpha+1)}$
Where

$$
\begin{aligned}
f_{1}(x)= & f(x) f^{\prime}(x)-f^{\prime \prime \prime}(x)+f^{\prime \prime}(x)-f^{\prime}(x) \\
& =(-\sin x-1) \cos x-2 \sin x \\
u_{2}= & -J^{\alpha}\left(A_{1}\right)+J^{\alpha}\left(L_{x x x} u_{1}\right)-J^{\alpha}\left(L_{x x} u_{1}\right)+J^{\alpha}\left(L_{x} u_{1}\right) \\
= & \frac{f_{2}(x)}{\Gamma(2 \alpha+1)} t^{2 \alpha}
\end{aligned}
$$

Where
$f_{2}(x)=-\left[f(x)+f_{1}^{\prime}(x)+f_{1}(x) f^{\prime}(x)-f_{1}^{\prime \prime \prime}(x)+\right.$ $\left.f_{1}^{\prime \prime}(x)-f_{1}^{\prime}(x)\right]$
$=\cos ^{3} x+3 \cos ^{2} x+((-\sin x-1) \sin x-2 \sin x-1) \cos x-5 \sin ^{2} x-2 \sin x$, similarly
$u_{3}=\frac{f_{3}(x)}{\Gamma(3 \alpha+1)} t^{3 \alpha}$
!
$u_{n}=\frac{f_{n}(x)}{\Gamma(n \alpha+1)} t^{n \alpha}$
$u(x, t)=f(x)+\frac{t^{\alpha}}{\Gamma(a+1)} f_{1}(x)+\frac{t^{2 \alpha}}{\Gamma(2 a+1)} f_{2}(x)+\cdots+\frac{t^{n a}}{\Gamma(n a+1)} f_{n}(x)+\cdots$
$=\sum_{n=0}^{\infty} \frac{t^{n \alpha}}{\Gamma(n \alpha+1)} f_{n}(x)$ Where $f_{0}(x)$ is an initial condition.

In the following example, we try to find the solve of another nonlinear fractional equation

## Example 3.

Consider the following of nonlinear fractional equation:-
$D_{t}^{\alpha} u+D_{x}^{2} u-D_{x} u+u^{2}=0$
$0<x \leq 1,0 \leq t \leq 1, \quad 0<\alpha \leq 1$
With initial condition

$$
\begin{aligned}
u(x, 0)=\varphi= & f(x)=x^{2} \\
& (x, t) \in \Omega \times[0, T]
\end{aligned}
$$

Note that here $\Omega=(0,1)$
The standard form of the fractional equation an operator form is
$D_{t}^{\alpha} u=-[u(x, t)]^{2}-L_{x x} u(x, t)+L_{x} u(x, t)$
Where $L_{x x}=\frac{\partial^{2}}{\partial x^{2}}, L_{x}=\frac{\partial}{\partial x}$ and the fractional differential operator $D_{t}^{\alpha}$ defined in (2) respectively.
$J^{\alpha}$ is the inverse of $D_{t}^{\alpha}$
Now applying ${ }^{\alpha}$ to the both side of our Eq. we obtain
$u(x, t)=\phi-J^{\alpha}(N u)-J^{\alpha}\left(L_{x x} u(x, t)\right)+J^{\alpha}\left(L_{x} u(x, t)\right)$
Where $N u=u^{2}$, according to the decomposition method, we assume series solution for the unknown function $u(x, t)$ in the form
$u(x, t)=\sum_{n=0}^{\infty} u_{n}(x, t)$
In order to solve our problem, we must generalize these Adomian polynomials as follows:

$$
\begin{array}{r}
A n=\frac{1}{n!}\left[\frac{d^{n}}{d \lambda^{n}}\left(\sum_{i=0}^{\infty} \lambda^{i} u_{i}\right)\left(\sum_{i=0}^{\infty} \lambda^{i} u_{i}\right)\right]_{\lambda=0} \\
, \quad n=0,1, \ldots
\end{array}
$$

$$
\begin{aligned}
& A_{0}=u_{0} \cdot u_{0}=u_{0}^{2} \\
& A_{1}=\frac{d}{d \lambda}\left[\left(u_{0}+\lambda u_{1}\right) \cdot\left(u_{0}+\lambda u_{1}\right)\right]_{\lambda=0} \\
& =u_{0} u_{1}+u_{1} u_{0}=2 u_{0} u_{1} \\
& A_{2}=\frac{1}{2!} \frac{d^{2}}{d \lambda^{2}}\left[\left(u_{0}+\lambda u_{1}+\lambda^{2} u_{2}\right)\left(u_{0}+\lambda u_{1}+\lambda^{2} u_{2}\right)\right] \\
& =2 u_{0} u_{2}+u_{1}^{2} \\
& A_{3}=\frac{1}{3!} \frac{d^{3}}{d \lambda^{3}}\left[\left(u_{0}+\lambda u_{1}+\lambda^{2} u_{2}+\lambda^{3} u_{3}\right)\left(u_{0}+\lambda u_{1}+\lambda^{2} u_{2}+\lambda^{3} u_{3}\right)\right]_{\lambda=0} \\
& =2 u_{0} u_{3}+2 u_{1} u_{2}
\end{aligned}
$$

;
And so on
From (7), we get
$\sum_{n=0}^{\infty} u_{n}(x, t)=\varphi-J^{a}\left(\sum_{n=0}^{\infty} A_{n}\left(u_{0}, u_{1}, \ldots, u_{n}\right)\right)-J^{\alpha}\left(\sum_{n=0}^{\infty}\left(L_{x x} u_{n}\right)\right)+J^{x}\left(\sum_{n=0}^{\infty}\left(L_{x} u_{n}\right)\right)$
From the above equation, we observe that
$u_{0}=u(x, 0)=f(x)$
$u_{1}=-J^{\alpha}\left(A_{0}\right)-J^{\alpha}\left(L_{x x} u_{0}\right)+J^{\alpha}\left(L_{x} u_{0}\right)$
$u_{2}=-J^{\alpha}\left(A_{1}\right)-J^{\alpha}\left(L_{x x} u_{1}\right)+J^{\alpha}\left(L_{x} u_{1}\right)$
!
$u_{n+1}=-J^{\alpha}\left(A_{n}\right)-J^{\alpha}\left(L_{x x} u_{n}\right)+J^{\alpha}\left(L_{x} u_{n}\right)$
By substituting the values of $u_{0}, u_{1}, \ldots$,
we get solutions of the IVP

$$
\begin{aligned}
& u=u_{0}+u_{1}+u_{2}+\ldots+u_{n}+\ldots \\
& u_{0}=u(x, 0)=f(x)=x^{2} \\
& \begin{aligned}
u_{1} & =-J^{\alpha}\left(A_{0}\right)-J^{\alpha}\left(L_{x x} u_{0}\right)+J^{\alpha}\left(L_{x}\left(u_{0}\right)\right) \\
& =\frac{t^{\alpha}}{\Gamma(\alpha+1)} f_{1}(x)
\end{aligned}
\end{aligned}
$$

Where
$f_{1}(x)=(f(x))^{2}+f^{\prime \prime}(x)-f^{\prime}(x)=x^{4}-2 x+2$
$\therefore u_{1}=x^{4}-2 x+2\left[\frac{t^{\alpha}}{\Gamma(\alpha+1)}\right]$
$u_{2}=-J^{\alpha}\left(A_{1}\right)-J^{\alpha}\left(L_{x x} u_{1}\right)+J^{\alpha}\left(L_{x} u_{1}\right)$
$=\frac{t^{2 \alpha}}{\Gamma(2 \alpha+1)} f_{2}(x)$
Where $f_{2}(x)=-\left[2 f(x) f_{1}(x)+f_{1}^{\prime \prime}(x)-f_{1}^{\prime}(x)\right]$

$$
=-2 x^{6}+8 x^{2}+2
$$

$u_{3}=-J^{\alpha}\left(A_{2}\right)-J^{\alpha}\left(L_{x x} u_{2}\right)+J^{\alpha}\left(L_{x} u_{2}\right)$
$=\frac{t^{3 \alpha}}{\Gamma(3 \alpha+1)} f_{3}(x)$
Where
$f_{3}(x)=-2 f(x) f_{2}(x)-2 f_{1}^{2}(x)-f_{2}^{\prime \prime}(x)+f_{2}^{\prime}(x)$

$$
=3 x^{8}-8 x^{5}+40 x^{4}-8 x^{2}+24 x-20
$$

Similarly
$u_{4}=\frac{t^{4 \alpha}}{\Gamma(4 \alpha+1)} f_{4}(x)$
:
$u_{n}=\frac{t^{n \alpha}}{\Gamma(n \alpha+1)} f_{n}(x)$
The solution of the considered IVP is given by
$u(x, t)=f(x)+\frac{t^{\alpha}}{\Gamma(\alpha+1)} f_{1}(x)+\frac{t^{2 \alpha}}{\Gamma(2 \alpha+1)} f_{2}(x)+\cdots$
$=\sum_{n=0}^{\infty} \frac{t^{n \alpha}}{\Gamma(n \alpha+1)} f_{n}(x)$

Where $f_{0}(x)=u(x, 0)$, is initial condition.

Also, in the next example, we solve a system of nonlinear equations with fractional orders

## Example 4.

Consider the system of initial value problem (IVP) of fractional equations
$D_{t}^{\alpha} u=u D_{x} u+v D_{y} u$
$D_{t}^{\alpha} v=u D_{x} v+v D_{y} v$
Where $0<\alpha \leq 1 \operatorname{and}(x, t) \in \Omega_{x}(0, T]$, and with the initial condition.
$u(x, y, 0)=f(x, y)$
$v(x, y, 0)=g(x, y) \quad, x, y \in \Omega$
Not that $\Omega=(0,1)$
The above system can be written in the equivalent form
$\frac{\partial^{\alpha} u}{\partial t^{\alpha}}=u \frac{\partial u}{\partial x}+v \frac{\partial u}{\partial y}$
$\frac{\partial^{\alpha} v}{\partial t^{\alpha}}=u \frac{\partial v}{\partial x}+v \frac{\partial v}{\partial y}$
$N_{1}(u, v)=u \frac{\partial u}{\partial x}+v \frac{\partial u}{\partial y}$
$N_{2}(u, v)=u \frac{\partial v}{\partial x}+v \frac{\partial v}{\partial y}$
$L u=N_{1}(u, v)$
$L v=N_{2}(u, v)$
Applying $L^{-1}()=.J^{\alpha}$ to both sides ofEq.(28) yields.
$u(x, y, t)=\phi+J^{\alpha} N_{1}(u, v)$
$v(x, y, t)=\phi+J^{\alpha} N_{2}(u, v)$
Where the nonlinear operator $N_{1}(u, v)$ and $N_{2}(u, v)$ are them written in the decomposition form
$N_{1}(u, v)=\sum_{n=0}^{\infty} A_{n}\left(u_{0}, u_{1}, \ldots, u_{n}\right)$
$N_{2}(u, v)=\sum_{n=0}^{\infty} B_{n}\left(v_{0}, v_{1}, \ldots, v_{n}\right)$

Where $A_{n}$ and $B_{n}$ are the Adomian polynomials of the following form

$$
\begin{array}{r}
A_{n}=\frac{1}{n!}\left[\frac{d^{n}}{d \lambda^{n}} N_{1}(u, v)\left(\sum_{i=0}^{\infty} \lambda^{i} u_{i}\right)\right]_{\lambda=0} \\
B_{n}=\frac{1}{n!}\left[\frac{d^{n}}{d \lambda^{n}} N_{2}(u, v)\left(\sum_{i=0}^{\infty} \lambda^{i} v_{i}\right)\right]_{\lambda=0} \\
, n=0,1, \ldots
\end{array}
$$

generalize these Adomian polynomials in follows.

$$
\begin{aligned}
& A_{n}=\frac{1}{n!}\left[\frac{d^{n}}{d \lambda^{n}}\left(\sum_{i=0}^{\infty} \lambda^{i} u_{i}\right)\left(\frac{\partial}{\partial x} \sum_{i=0}^{\infty} \lambda^{i} u_{i}\right)+\right. \\
& \left.\quad\left(\sum_{i=0}^{\infty} \lambda^{i} v_{i}\right)\left(\frac{\partial}{\partial y} \sum_{i=0}^{\infty} \lambda^{i} u_{i}\right)\right]_{\lambda=0}, \\
& n=0,1, \ldots \\
& A_{0}=u_{0} \frac{\partial u_{0}}{\partial x}+v_{0} \frac{\partial u_{0}}{\partial y} \\
& \therefore A_{1}=u_{0} \frac{\partial u_{1}}{\partial x}+v_{0} \frac{\partial u_{1}}{\partial y}+u_{1} \frac{\partial u_{0}}{\partial x}+v_{1} \frac{\partial u_{0}}{\partial y} \\
& \therefore A_{2}=u_{0} \frac{\partial u_{2}}{\partial x}+v_{0} \frac{\partial u_{2}}{\partial y}+u_{1} \frac{\partial u_{1}}{\partial x}+v_{1} \frac{\partial u_{1}}{\partial y}+u_{2} \frac{\partial u_{0}}{\partial x}+v_{2} \frac{\partial u_{0}}{\partial y}
\end{aligned}
$$

Similarly

$$
\therefore A_{3}=u_{0} \frac{\partial u_{3}}{\partial x}+v_{0} \frac{\partial u_{3}}{\partial y}+u_{1} \frac{\partial u_{2}}{\partial x}+v_{1} \frac{\partial u_{2}}{\partial y}+u_{2} \frac{\partial u_{1}}{\partial x}+v_{2} \frac{\partial u_{1}}{\partial y}+u_{3} \frac{\partial u_{0}}{\partial x}+v_{3} \frac{\partial u_{0}}{\partial y}
$$

$$
\vdots
$$

And so on
Now, to calculated our problem we must generalize these Adomian polynomials in follows
$B_{n}=\frac{1}{n!}\left[\frac{d^{n}}{d \lambda^{n}}\left(\sum_{i=0}^{\infty} \lambda^{i} u_{i}\right)\left(\frac{\partial}{\partial x} \sum_{i=0}^{\infty} \lambda^{i} v_{i}\right)+\left(\sum_{i=0}^{\infty} \lambda^{i} v_{i}\right)\left(\frac{\partial}{\partial y} \sum_{i=0}^{\infty} \lambda^{i} v_{i}\right)\right]_{\lambda=0} \quad \begin{aligned} & u_{2}=J^{\alpha} A_{1} \\ & v_{2}=J^{\alpha} B_{1}\end{aligned}$

$$
n=0,1,2, \ldots
$$

$u_{n+1}=J^{\alpha} A_{n}$
$v_{n+1}=J^{\alpha} B_{n}$
$B_{0}=u_{0} \frac{\partial u_{0}}{\partial x}+v_{0} \frac{\partial v_{0}}{\partial y}$
By substituting the values of $u_{0}, u_{1}, \ldots$ And $v_{0}, v_{1}, \ldots$ we get a solution of the
$B_{1}=u_{0} \frac{\partial v_{1}}{\partial x}+v_{0} \frac{\partial v_{1}}{\partial y}+u_{1} \frac{\partial v_{0}}{\partial x}+v_{1} \frac{\partial v_{0}}{\partial y}$ IVP.
$B_{2}=u_{0} \frac{\partial v_{2}}{\partial x}+v_{0} \frac{\partial v_{2}}{\partial y}+u_{1} \frac{\partial v_{1}}{\partial x}+v_{1} \frac{\partial v_{1}}{\partial y}+u_{2} \frac{\partial v_{0}}{\partial x}+v_{2} \frac{\partial v_{0}}{\partial y}$

$$
\therefore B_{3}=u_{0} \frac{\partial v_{3}}{\partial x}+v_{0} \frac{\partial v_{3}}{\partial y}+u_{1} \frac{\partial v_{2}}{\partial x}+v_{1} \frac{\partial v_{2}}{\partial y}+u_{2} \frac{\partial v_{1}}{\partial x}+v_{2} \frac{\partial v_{1}}{\partial y}+u_{3} \frac{\partial v_{0}}{\partial x}+v_{3} \frac{\partial v_{0}}{\partial y}
$$

$$
\begin{aligned}
& u_{0}=u(x, y, 0)=f(x, y) \\
& v_{0}=v(x, y, 0)=g(x, y) \\
& u_{1}=J^{\alpha}\left(A_{0}\right)=\frac{1}{\Gamma(\alpha)} \int_{0}^{t} \frac{1}{(t-\theta)^{1-\alpha}}\left[u_{0} \frac{\partial u_{0}}{\partial x}+v_{0} \frac{\partial u_{0}}{\partial y}\right] d \theta
\end{aligned}
$$

And so on
From (7), we obtain
$\sum_{n=0}^{\infty} u_{n}(\mathrm{x}, \mathrm{y}, \mathrm{t})=\mathrm{u}(\mathrm{x}, \mathrm{y}, 0)+$

$$
+I^{\alpha}\left(\sum_{n=0}^{\infty} A_{n}\left(u_{0}, u_{1}, \ldots, u_{n}\right)\right)
$$

$\sum_{n=0}^{\infty} v_{n}(\mathrm{x}, \mathrm{y}, \mathrm{t})=v(\mathrm{x}, \mathrm{y}, 0)+$

$$
+J^{\alpha}\left(\sum_{n=0}^{\infty} B_{n}\left(u_{0}, u_{1}, \ldots, u_{n}\right)\right)
$$

The associated decomposition is given by
$u_{0}=u(x, y, 0), u_{n+1}=I^{\alpha}\left(N_{1}\left(u_{n}, v_{n}\right)\right)$
$v_{0}=v(x, y, 0), v_{n+1}=I^{\alpha}\left(N_{2}\left(u_{n}, v_{n}\right)\right), \mathrm{n}=0,1,2, \ldots$

$$
u_{2}=J^{\alpha}\left(A_{1}\right)=f_{2}(x) \frac{t^{2 \alpha}}{\Gamma(2 \alpha+1)}
$$

Then, According to the above equations we get,
$u_{0}=u(x, y, 0)$
$v_{0}=v(x, y, 0)$
$u_{1}=J^{\alpha} A_{0}$
$v_{1}=J^{\alpha} B_{0}$
Where

$$
\begin{aligned}
& f_{1}(x, y)=-\left[f(x, y) \frac{\partial f(x, y)}{\partial x}+g(x, y) \frac{\partial f(x, y)}{\partial y}\right] \\
& v_{1}=J^{\alpha} B_{0}= \\
& =\frac{1}{\Gamma(\alpha)} \int_{0}^{t} \frac{1}{(t-\theta)^{1-\alpha}}\left[u_{0} \frac{\partial v_{0}}{\partial x}+v_{0} \frac{\partial v_{0}}{\partial y}\right] d \theta \\
& \\
& =g_{1}(x, y) \frac{t^{\alpha}}{\Gamma(\alpha+1)}
\end{aligned}
$$

Where

$$
g_{1}(x, y)=-\left[f(x, y) \frac{\partial g(x, y)}{\partial x}+g(x, y) \frac{\partial g(x, y)}{\partial y}\right]
$$

Where

$$
\begin{aligned}
& f_{2}(x)=\left[f(x, y) \frac{\partial f_{1}(x, y)}{\partial x}+g(x, y) \frac{\partial f_{1}(x, y)}{\partial y}+f_{1}(x, y) \frac{f(x, y)}{\partial x}+g_{1}(x, y) \frac{\partial f(x, y)}{\partial y}\right] \\
& v_{2}=J^{\alpha} B_{1}=g_{2}(x, y) \frac{t^{2 \alpha}}{\Gamma(2 \alpha+1)}
\end{aligned}
$$

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Where
$g_{2}(x, y)=f(x, y) \frac{\partial g_{1}(x, y)}{\partial x}+g(x, y) \frac{\partial g_{1}(x, y)}{\partial y}+$
$f_{1}(x, y) \frac{\partial g(x, y)}{\partial x}+g_{1}(x, y) \frac{\partial g(x, y)}{\partial y}$

By induction, we have
$u(x, y, t)=\sum_{n=0}^{\infty} u_{n}=u_{0}+u_{1}+u_{2}+\cdots+u_{n}+\cdots$
$=f(x, y)+f_{1}(x, y) \frac{t^{\alpha}}{\Gamma(\alpha+1)}+f_{2}(x, y) \frac{t^{2 \alpha}}{\Gamma(2 \alpha+1)}+$
$v(x, y, t)=\sum_{n=0}^{\infty} v_{n}=v_{0}+v_{1}+\cdots+v_{n}+\cdots$
$=g(x, y)+g_{1}(x, y) \frac{t^{a}}{\Gamma(a+1)}+g_{2}(x, y) \frac{t^{2 a}}{\Gamma(2 a+1)}+\cdots+g_{n}(x, y) \frac{t^{n a}}{\Gamma(n a+1)}+\cdots$

Finally, we apply the Adomian decomposition method to obtain approximate solutions of IVPs for fractional BBM-Burger's equation with $\varepsilon=1$

## Example 5.

Consider the initial value problem (IVP) for fractional BBM_Burger's equation of the form
$\frac{\partial^{\alpha} u}{\partial t^{\alpha}}-\frac{\partial^{2} u}{\partial x^{2}}+u \frac{\partial u}{\partial x}=0$
Where $0<\alpha \leq 1$ and with initial condition
$u(x, 0)=\varphi=f(x)=\sin (x), x \in \Omega \times(0, T]$
Note that here $\Omega=(0,1)$, the standard form of the fractional BBM_Burger's equation in an operator form
$\frac{\partial^{\alpha} u}{\partial t^{\alpha}}=\frac{\partial^{2} u}{\partial x^{2}}-u \frac{\partial u}{\partial x}$
$\frac{\partial^{\alpha} u(x, t)}{\partial t^{\alpha}}=\left(L_{x x} u(x, t)\right)-\left(u(x, t) L_{x} u(x, t)\right)$

Where $L_{x x}=\frac{\partial^{2}}{\partial x^{2}}, \quad L_{x}=\frac{\partial}{\partial x}$
And the fractional differential operator $\frac{\partial^{\alpha}}{\partial t^{\alpha}}$ defined in equation (2), respectively we know that $J^{\alpha}$ which is invers of the operator $\frac{\partial^{\alpha}}{\partial t^{\alpha}}$, now applying $J^{\alpha}$ to the both sides of our Eq., we get
$u(x, t)=\phi+J^{\alpha}\left(L_{x x} u\right)-J^{\alpha}(N u)$

In order to solve our problem we must generalize these Adomian polynomials as follows.
$A_{n}=\frac{1}{n!} \frac{d^{n}}{d \lambda^{n}}\left[\left(\sum_{i=0}^{n} \lambda^{i} u_{i}\right) \frac{\partial}{\partial x}\left(\sum_{i=0}^{n} \lambda^{i} u_{i}\right)\right]_{\lambda=0}$
$n=0,1,2, \ldots$
$A_{0}=u_{0} \frac{\partial u_{0}}{\partial x}$
$A_{1}=u_{0} \frac{\partial u_{1}}{\partial x}+u_{1} \frac{\partial u_{0}}{\partial x}$
$A_{2}=u_{0} \frac{\partial u_{2}}{\partial x}+u_{1} \frac{\partial u_{1}}{\partial x}+u_{2} \frac{\partial u_{0}}{\partial x}$
:
And so on
Thus
$u(x, t)=\varphi+J^{\alpha}\left(\sum_{n=0}^{\infty}\left(L_{x x} u_{n}\right)\right)-J^{\alpha}\left(\sum_{n=0}^{\infty}\left(A_{n}\right)\right)$,
$u_{0}=\varphi=u(x, 0)$
$u_{1}=J^{\alpha}\left(L_{x x} u_{0}\right)-J^{\alpha}\left(A_{0}\right)$
$u_{2}=J^{\alpha}\left(L_{x x} u_{1}\right)-J^{\alpha}\left(A_{1}\right)$
!
$u_{n+1}=J^{\alpha}\left(L_{x x} u_{n}\right)-J^{\alpha}\left(A_{n}\right)$
Consequently
$u(x, t)=u_{0}+u_{1}+u_{2}+\cdots+u_{n}+\cdots$
$u_{0}=u(x, 0)=f(x)=\sin x$
$u_{1}=J^{\alpha}\left(L_{x x} u_{0}\right)-J^{\alpha}\left(A_{0}\right)=\frac{f_{1}(x)}{\Gamma(\alpha+1)} t^{\alpha}$
Where $f_{1}(x)=-f^{\prime \prime}(x)+f(x) f^{\prime}(x)$
$=\sin x)(1+\cos x)$
$u_{2}=J^{\alpha}\left(L_{x x} u_{1}\right)-J^{\alpha}\left(A_{1}\right)=\frac{f_{2}(x)}{\Gamma(\alpha+1)} t^{\alpha}$
Where
$f_{2}(x)=$
$f_{1}^{\prime \prime}(x)-f(x) f_{1}^{\prime}(x)-$ $f_{1}(x) f^{\prime}(x)$
nonlinear partial equations. We demonstrated that the decomposition procedure is quite efficient to determine the exact solutions. However, the method gives a simple powerful tool for obtaining the solutions without a need for large size of computations. It is also worth noting that the advantage of this method sometimes displays a fast convergence of the solutions. In addition, the numerical results which obtained by this method indicate a high degree of accuracy, also efficiency of the desired results.

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$u_{3}=J^{\alpha}\left(L_{x x} u_{2}\right)-J^{\alpha}\left(A_{2}\right)=f_{3}(x) \frac{t^{3 \alpha}}{\Gamma(3 \alpha+1)}$
Similarly
$\therefore u_{3}=\frac{f_{3}(x)}{\Gamma(3 \alpha+1)} t^{3 \alpha}$
!
$u_{n}=\frac{f_{n}(x)}{\Gamma(n \alpha+1)} t^{n \alpha}$
The solution of the considered IVP is given by
$u(x, t)=f(x)+\frac{t^{\alpha}}{\Gamma(\alpha+1)} f_{1}(x)+\frac{t^{2 \alpha}}{\Gamma(2 \alpha+1)} f_{2}(x)+$
$\frac{t^{3 \alpha}}{\Gamma(3 \alpha+1)} f_{3}(x)+\cdots+\frac{t^{n \alpha}}{\Gamma(n \alpha+1)} f_{n}(x)+\cdots$
(38)
$=\sum_{n=0}^{\infty} \frac{t^{n \alpha}}{\Gamma(n \alpha+1)} f_{n}(x)$,
where $f_{0}(x)$ is an initial condition. (Comp. 7-23).

## Conclusion

In this paper, we have applied the Adomian decomposition method for solving problems of
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